# Quantized Tensor Trains for Integral Equations 

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$$
\begin{gathered}
a \sigma(x)+\int_{\Omega} K(x, y) \sigma(y) d \Omega_{y}=f(x) \\
A \sigma=f
\end{gathered}
$$

Uses:

- Various physical phenomena (especially $K(x, y)=\kappa(|x-y|)$ ).
- Mesh-free methods for PDEs.


## Quick Reminder: Solving Matrix Equations

GMRES, CG, et al. minimize a residual norm $\sigma \mapsto\|A \sigma-f\|$.
To compute this, we only need a map $\sigma \mapsto A \sigma$.
If $A$ is poorly-conditioned, this can converge slowly.
If we can define $P$ such that $P(x) \approx A^{-1} x$, we can solve $P(A) \sigma=P(f)$ instead.

## The Plan

Idea: Use a QTT to represent $A$.

1. Factor $A$ into a QTT.
2. Do matvecs by QTTs (which lets us invert $A$ with GMRES).
3. Do approximate inverses in QTT form (which lets us precondition GMRES).

## Step 1: Quantizing $A$

Idea: Split $A$ into a grid which defines a source tree and a target tree.
Definition: A box in the source tree and a box in the target tree are well-separated if the distance between them is at least the diameter of the largest box.

Intuition: Well-separated source-leaf/target-leaf pairs are far from the main diagonal of the matrix $A$.

Definition: $A$ is FMM-compressible if any matrix sub-block representing a wellseparated pair of boxes has maximum numerical rank $k_{\epsilon}$ for some fixed $k_{\epsilon}$.

## Tensorizing $A$

Form the product tree, where every node represents a node in the source and the target trees at once.

Think of $A$ as a tensor where each element is given by a path through the product tree.

Note that unfolding matrices of $A$ are not the same as the original matrix!

## Does This Work?

Definition: A kernel $K(x, y)$ is trans/ation-invariant if $K(x+\omega, y+\omega)=K(x, y)$ for all $\omega$.

Theorem: Let $A$ be FMM-compressible and let the kernel $K$ be translationinvariant. Then the QTT representation has bounded rank $r=\max \left(r_{k}\right) \leq k_{\epsilon}^{2}+4 D-1$, where $D$ is the dimension of the ambient space.

## Step 2: Matvecs by QTT-Compressed Matrices

Let $A$ have cores given by $G_{k}^{A}\left(\alpha_{k-1}, \overline{i_{k} j_{k}}, \alpha_{k}\right)$ and $b$ have cores given by $G_{k}^{b}\left(\beta_{k-1}, j_{k}, \beta_{k}\right)$. Then $A b$ has cores given by

$$
G_{k}^{A b}\left(\overline{\alpha_{k-1} \beta_{k-1}}, i_{k}, \overline{\alpha_{k} \beta_{k}}\right)=\sum_{j_{k}} G_{k}^{A}\left(\alpha_{k-1} \overline{i_{k} j_{k}} \alpha_{k}\right) G_{k}^{b}\left(\beta_{k-1} j_{k} \beta_{k}\right) .
$$

$O\left(r_{A}^{2} r_{b}^{2} N \log N\right)$.
Well, that was easy.

## Step 2, Shia Surprise: Matvecs with Uncompressed Vectors

Idea: Split $b$ up along the target tree and use QTT cores to move it to the source tree.

Initialize: $y_{0}=b^{T}=y_{0}\left(\alpha_{0}, \overline{j_{1} j_{2} \cdots j_{d}}\right)$
Core iteration:

1. $M_{k}\left(\overline{\alpha_{k} i_{k}}, \overline{\alpha_{k-1} j_{k}}\right)=G_{k}^{A}\left(\alpha_{k-1}, \overline{i_{k} j_{k}}, \alpha_{k}\right)$
2. $b_{k}\left(\overline{\alpha_{k-1} j_{k}}, \overline{i_{1} \cdots i_{k-1} j_{k+1} \cdots j_{d}}\right)=y_{k-1}\left(\alpha_{k-1}, \overline{i_{1} \cdots i_{k-1} j_{k} \cdots j_{d}}\right)$
3. $\phi_{k}=M_{k} b_{k}$
4. $y_{k}\left(\alpha_{k}, \overline{i_{1} \cdots i_{k} j_{k+1} \cdots j_{d}}\right)=\phi_{k}\left(\overline{\alpha_{k} i_{k}}, \overline{i_{1} \cdots i_{k-1} j_{k+1} \cdots j_{d}}\right)$
$O\left(r_{A}^{2} N \log N\right)$.

## Step 3: (Approximate) Inversion of QTTCompressed Matrices

Idea: Consider $A X=I$ and solve for $X$.
Useful identity: $\operatorname{vec}(A B C)=\left(C^{T} \otimes A\right) \operatorname{vec}(B) \Longrightarrow(I \otimes A) \operatorname{vec}(X)=\operatorname{vec}(I)$
Write $X$ as a QTT, fix all cores but one, and solve for that core. Repeat for all cores.
Problem: can't change the ranks of $X$.
One possible solution (DMRG): solve for two cores at once, then split them up.

## Reference

Corona, E., Rahimian, A., \& Zorin, D. (2017). A Tensor-Train accelerated solver for integral equations in complex geometries. Journal of Computational Physics, 334, 145-169. https://doi.org/10.1016/j.jcp.2016.12.051

